SCATTERING IN CHERN-SIMONS MATTER THEORIES

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INTRODUCTION

CHERN-SIMONS THEORY: WHAT

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 $\mathcal{L}_{\mathrm{CS}}$ is NOT gauge invariant. But if the **level** $k \in \mathbb{Z}$ then we note that

 $\exp(iS_{\rm CS}[A])$

is unchanged. This is all we care about for our quantum theory.

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- Canonical example of a topological quantum field theory
- Knot invariants are expectation values of Wilson loops in Chern-Simons theories
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Some neat facts about Chern-Simons theories:

- $\circ~$ Time reversal invariant depending on level of k and gauge group
- Bosonization duality between Chern-Simons with bosons and with fermions
- \circ Large N Chern-Simons matter theories are dual to higher spin gravity in AdS_4

CHERN-SIMONS MATTER THEORIES

Chern-Simons with U(N) gauge fields coupled to fundamental bosons and fermions

$$\mathcal{L}_{\text{Bos}} = \mathcal{L}_{\text{CS}} + D_{\mu}\bar{\phi}D^{\mu}\phi + m_{B}^{2}\bar{\phi}\phi + \frac{b_{4}}{2N_{b}}(\bar{\phi}\phi)^{2}$$
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Supported by evidence:

- 3pt functions match in $N \to \infty$
- Thermal partition functions matching in $N \rightarrow \infty$
- SUSY duality can be constructed which flows to these ones after SUSY breaking

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Causality is usually upgraded to a slightly stronger condition: that transition amplitudes are real-boundary values of analytic functions F(s, t, u) in complex *s*-plane.

$$S = \mathbb{1} + i (2\pi)^d \,\delta^{(d)}(p_1 + p_2 - p_3 - p_4) \,F(s, t, u)$$

S-MATRIX: CROSSING-SYMMETRY



Crossing symmetry means the amplitude for something like

 $s: A_1 + A_2 \to A_3 + A_4$

can be obtained from

 $u: A_1 + \bar{A_4} \to A_3 + \bar{A_2}$

by analytic continuation of F(s, t, u).

THE PROJECT

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- Believe this factor comes from a need to attach Wilson lines to fermions to make correlation functions gauge invariant.
- Not clear if what they found is a large N artifact. Want to do calculation in usual perturbative regime $(k \to \infty)$ but N finite.

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Photon Correction







- There are two diagrams that have an anomalous magnetic moment-type correction
- There is also a crossed box.

Expressions are quite large and need to be tackled algebraically



$$= -e^{4} \epsilon_{\mu\nu\lambda} \epsilon_{\sigma\rho\beta} \int \frac{d^{3}k}{(2\pi)^{3}} \frac{(\bar{u}_{3}\gamma^{\rho} [\not{l} + M]\gamma^{\mu} u_{1})}{l^{2} - M^{2}} \\ \times \frac{(\bar{v}_{2}\gamma^{\sigma} [\not{p}_{1} - \not{l} - \not{p}_{4} + M]\gamma^{\nu} v_{4})}{(p_{1} - l - p_{4})^{2} - M^{2}} \\ \times \frac{(p_{1} - l)^{\lambda} (p_{3} - k)^{\beta}}{(p_{1} - l)^{2} (p_{3} - l)^{2}}$$

1. Take diagram input and reduce products of gamma functions using

$$\gamma^{\alpha}\gamma^{\beta} = g^{\alpha\beta} - i\epsilon^{\alpha\beta\sigma}\gamma_{\sigma}$$

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4. Collect the expression into sums of terms proportional to known master integrals

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3. For the box diagrams we cannot solve the integrals with Feynman parameters, we trade them for a (double) Mellin-Barnes integral

$$\frac{1}{(X+Y)^{\lambda}} = \frac{1}{\Gamma(\lambda)} \frac{1}{2\pi i} \int_{-i\infty}^{+i\infty} dz \, \Gamma(\lambda+z) \Gamma(-z) \frac{Y^z}{X^{\lambda+z}}$$

OUR RESEARCH: ANALYTIC PROBLEMS

- 4. Completed computation of anomalous magnetic moment and bubble diagram.
- 5. Checked that these diagrams satisfy the Ward identities imposed by gauge invariance (which had hundreds of terms).



 q^{μ} 0 \times _

CONCLUSION

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FUTURE DIRECTIONS:

- Finish the box integrals
- Combine all the expressions
- Test crossing symmetry

