

HOLOMORPHIC QFTs: HIGHER STRUCTURES AND BOOTSTRAP

STRINGS 2022

JUSTIN KULP

WITH KASIA BUDZIK, DAVIDE GAIOTTO, BRIAN WILLIAMS,
JINGXIANG WU AND MATTHEW YU.

PERIMETER INSTITUTE FOR
THEORETICAL PHYSICS

21/JUL/2022

HOLOMORPHIC QFTS AND HOLOMORPHIC TWISTS

- **Holomorphic QFTs** depend only on the complex structure of the underlying manifold [Johansen], [Nekrasov], [Costello], [Williams],
 - ▶ Local operators carry structure of a **Holomorphic Factorization Algebra** [Costello, Gwilliam].
 - ▶ Example. 2d Chiral Algebra and/or Vertex Algebra

HOLOMORPHIC QFTS AND HOLOMORPHIC TWISTS

- **Holomorphic QFTs** depend only on the complex structure of the underlying manifold [Johansen], [Nekrasov], [Costello], [Williams], ...
 - ▶ Local operators carry structure of a **Holomorphic Factorization Algebra** [Costello, Gwilliam].
 - ▶ Example. 2d Chiral Algebra and/or Vertex Algebra
- Given any SQFT, we obtain the **Holomorphic Twist** by taking cohomology of any one nilpotent supercharge, e.g. $Q := Q_-$
 - ▶ Anti-holomorphic translations are Q -exact, so twisted theory is (cohomologically) holomorphic

$$\{Q, \bar{Q}_{\dot{\alpha}}\} = \partial_{\bar{z}^{\dot{\alpha}}} \quad (1)$$

- ▶ Infinite dimensional symmetry enhancements analogous to Virasoro and Kac-Moody [Gwilliam, Williams].
- ▶ Deformation \rightsquigarrow [Beem, Lemos, Liendo, Peelaers, Raselli, van Rees]

HOLOMORPHIC QFTS AND HOLOMORPHIC TWISTS

- **Holomorphic QFTs** depend only on the complex structure of the underlying manifold [Johansen], [Nekrasov], [Costello], [Williams], ...
 - ▶ Local operators carry structure of a **Holomorphic Factorization Algebra** [Costello, Gwilliam].
 - ▶ Example. 2d Chiral Algebra and/or Vertex Algebra
- Given any SQFT, we obtain the **Holomorphic Twist** by taking cohomology of any one nilpotent supercharge, e.g. $Q := Q_-$
 - ▶ Anti-holomorphic translations are Q -exact, so twisted theory is (cohomologically) holomorphic

$$\{Q, \bar{Q}_{\dot{\alpha}}\} = \partial_{\bar{z}^{\dot{\alpha}}} \quad (1)$$

- ▶ Infinite dimensional symmetry enhancements analogous to Virasoro and Kac-Moody [Gwilliam, Williams].
 - ▶ Deformation \rightsquigarrow [Beem, Lemos, Liendo, Peelaers, Raselli, van Rees]
- Operators captured by holomorphic twist are those counted by **superconformal index** [Saber, Williams]

$$\mathcal{I} = \text{Tr}(-1)^F p^{j_1+j_2-r/2} q^{j_1-j_2-r/2} e^{-\beta\{Q_-, S^-\}} \quad (2)$$

HIGHER BRACKETS AND HOMOTOPY TRANSFER

- Theories are equipped with local product called **λ -Bracket**

$$\{\mathcal{O}_1, \mathcal{O}_2\}_\lambda = \oint_{S^3} e^{\lambda \cdot z} d^2 z \mathcal{O}_1(z, \bar{z}) \mathcal{O}_2(0) \quad (3)$$

- ▶ **Higher brackets** describe homotopy between lower brackets

$$\{\mathcal{O}_1, \mathcal{O}_2, \dots, \mathcal{O}_{n+1}\}_{\lambda_1, \dots, \lambda_n} \quad (4)$$

HIGHER BRACKETS AND HOMOTOPY TRANSFER

- Theories are equipped with local product called **λ -Bracket**

$$\{\mathcal{O}_1, \mathcal{O}_2\}_\lambda = \oint_{S^3} e^{\lambda \cdot z} d^2 z \mathcal{O}_1(z, \bar{z}) \mathcal{O}_2(0) \quad (3)$$

- ▶ **Higher brackets** describe homotopy between lower brackets

$$\{\mathcal{O}_1, \mathcal{O}_2, \dots, \mathcal{O}_{n+1}\}_{\lambda_1, \dots, \lambda_n} \quad (4)$$

- Polynomials in fields and derivatives \rightsquigarrow **Free Cohomology** \mathcal{V}

- ▶ **Interacting quantum theory** is obtained from underlying free-classical theory \mathcal{V} as cohomology of a new operator

$$\mathbf{Q} = Q_0 + Q_1 + Q_2 \dots \quad (5)$$

where Q_n is computed by n -loop Feynman diagrams.

HIGHER BRACKETS AND HOMOTOPY TRANSFER

- Theories are equipped with local product called **λ -Bracket**

$$\{\mathcal{O}_1, \mathcal{O}_2\}_\lambda = \oint_{S^3} e^{\lambda \cdot z} d^2 z \mathcal{O}_1(z, \bar{z}) \mathcal{O}_2(0) \quad (3)$$

- ▶ **Higher brackets** describe homotopy between lower brackets

$$\{\mathcal{O}_1, \mathcal{O}_2, \dots, \mathcal{O}_{n+1}\}_{\lambda_1, \dots, \lambda_n} \quad (4)$$

- Polynomials in fields and derivatives \rightsquigarrow **Free Cohomology** \mathcal{V}

- ▶ **Interacting quantum theory** is obtained from underlying free-classical theory \mathcal{V} as cohomology of a new operator

$$\mathbf{Q} = Q_0 + Q_1 + Q_2 \dots \quad (5)$$

where Q_n is computed by n -loop Feynman diagrams.

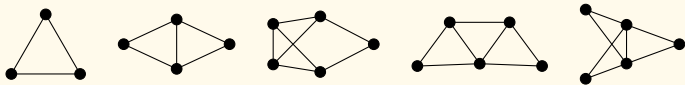
- **All perturbative corrections** are contained in the higher brackets of the free holomorphic factorization algebra!

$$\mathbf{Q}\mathcal{O} = \{\mathcal{I}, \mathcal{O}\}_0 + \{\mathcal{I}, \mathcal{I}, \mathcal{O}\}_{0,0} + \{\mathcal{I}, \mathcal{I}, \mathcal{I}, \mathcal{O}\}_{0,0,0} + \dots \quad (6)$$

- ▶ [Tree-level] $\sim 1 \mathcal{I}$, [1-Loop] $\sim 2 \mathcal{I}$'s, etc.

FEYNMAN DIAGRAMS AND BOOTSTRAP

- Feynman diagrams in theory must be **Laman graphs**.

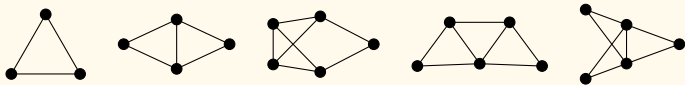


- ▶ Arbitrary integral takes the form:

$$\mathcal{I}_\Gamma[\lambda; z] \equiv \int_{\mathbb{R}^{4|\Gamma_0|-4}} \bar{\partial} \left[\prod_{e \in \Gamma_1} \mathcal{P}(x_{e_0} - x_{e_1} + z_e, \bar{x}_{e_0} - \bar{x}_{e_1}) \right] \left[\prod_{v \in \Gamma'_0} e^{\lambda_v \cdot x_v} d^2 x_v \right] \quad (7)$$

FEYNMAN DIAGRAMS AND BOOTSTRAP

- Feynman diagrams in theory must be **Laman graphs**.



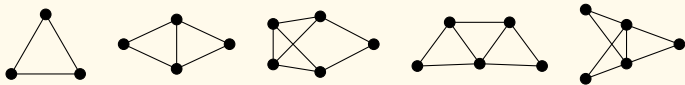
- ▶ Arbitrary integral takes the form:

$$\mathcal{I}_\Gamma[\lambda; z] \equiv \int_{\mathbb{R}^{4|\Gamma_0|-4}} \bar{\partial} \left[\prod_{e \in \Gamma_1} \mathcal{P}(x_{e_0} - x_{e_1} + z_e, \bar{x}_{e_0} - \bar{x}_{e_1}) \right] \left[\prod_{v \in \Gamma'_0} e^{\lambda_v \cdot x_v} d^2 x_v \right] \quad (7)$$

- Change of variables maps integral to Fourier transform of a polytope in space of holomorphic loop momenta.

FEYNMAN DIAGRAMS AND BOOTSTRAP

- Feynman diagrams in theory must be **Laman graphs**.



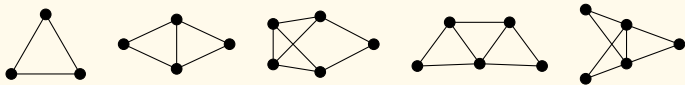
- ▶ Arbitrary integral takes the form:

$$\mathcal{I}_\Gamma[\lambda; z] \equiv \int_{\mathbb{R}^{4|\Gamma_0|-4}} \bar{\partial} \left[\prod_{e \in \Gamma_1} \mathcal{P}(x_{e_0} - x_{e_1} + z_e, \bar{x}_{e_0} - \bar{x}_{e_1}) \right] \left[\prod_{v \in \Gamma'_0} e^{\lambda_v \cdot x_v} d^2 x_v \right] \quad (7)$$

- Change of variables maps integral to Fourier transform of a polytope in space of holomorphic loop momenta. The **Holohedron?**

FEYNMAN DIAGRAMS AND BOOTSTRAP

- Feynman diagrams in theory must be **Laman graphs**.



- ▶ Arbitrary integral takes the form:

$$\mathcal{I}_\Gamma[\lambda; z] \equiv \int_{\mathbb{R}^{4|\Gamma_0|-4}} \bar{\partial} \left[\prod_{e \in \Gamma_1} \mathcal{P}(x_{e_0} - x_{e_1} + z_e, \bar{x}_{e_0} - \bar{x}_{e_1}) \right] \left[\prod_{v \in \Gamma'_0} e^{\lambda_v \cdot x_v} d^2 x_v \right] \quad (7)$$

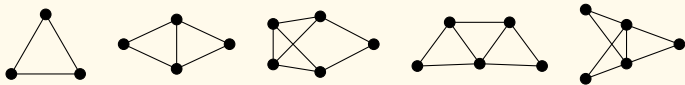
- Change of variables maps integral to Fourier transform of a polytope in space of holomorphic loop momenta. The **Holohedron**?

- ▶ **Configuration spaces** of graphs satisfy infinite collection of geometric **quadratic identities**; enforcing associativity

$$\sum_S \sigma(\Gamma, S) \mathcal{I}_{\Gamma[S]}[\lambda + \partial_{z'}; z] \mathcal{I}_{\Gamma(S)}[\lambda'; z'] = 0. \quad (8)$$

FEYNMAN DIAGRAMS AND BOOTSTRAP

- Feynman diagrams in theory must be **Laman graphs**.



- ▶ Arbitrary integral takes the form:

$$\mathcal{I}_\Gamma[\lambda; z] \equiv \int_{\mathbb{R}^{4|\Gamma_0|-4}} \bar{\partial} \left[\prod_{e \in \Gamma_1} \mathcal{P}(x_{e_0} - x_{e_1} + z_e, \bar{x}_{e_0} - \bar{x}_{e_1}) \right] \left[\prod_{v \in \Gamma'_0} e^{\lambda_v \cdot x_v} d^2 x_v \right] \quad (7)$$

- Change of variables maps integral to Fourier transform of a polytope in space of holomorphic loop momenta. The **Holohedron**?

- ▶ **Configuration spaces** of graphs satisfy infinite collection of geometric **quadratic identities**; enforcing associativity

$$\sum_S \sigma(\Gamma, S) \mathcal{I}_{\Gamma[S]}[\lambda + \partial_{z'}; z] \mathcal{I}_{\Gamma(S)}[\lambda'; z'] = 0. \quad (8)$$

- Find that quadratic identities are sufficient to **bootstrap Feynman integrals** to at least 3-loops (perhaps further!)

FIN