Holomorphic Confinement of $\mathcal{N}=1$ SYM PiTP 2023

JUSTIN KULP Kasia Budzik, Davide Gaiotto, Brian Williams, Jingxiang Wu, Matt Yu.

PERIMETER INSTITUTE FOR THEORETICAL PHYSICS

13/JUL/2023

ARXIV:2207.14321 ARXIV:2306.01039

HOLOMORPHIC TWISTS AND INFINITE SYMMETRIES

- Given any SQFT, we obtain the **Holomorphic Twist** by taking cohomology of any nilpotent supercharge, e.g. *Q* := *Q*_.
 - Anti-holomorphic translations are Q-exact, so twisted theory is (cohomologically) holomorphic [Johansen], [Witten], [Costello]

$$\{Q, \bar{Q}_{\dot{\alpha}}\} = \partial_{\bar{z}^{\dot{\alpha}}} \tag{1}$$

HOLOMORPHIC TWISTS AND INFINITE SYMMETRIES

- Given any SQFT, we obtain the **Holomorphic Twist** by taking cohomology of any nilpotent supercharge, e.g. *Q* := *Q*_.
 - Anti-holomorphic translations are Q-exact, so twisted theory is (cohomologically) holomorphic [Johansen], [Witten], [Costello]

$$\{Q, \bar{Q}_{\dot{\alpha}}\} = \partial_{\bar{z}^{\dot{\alpha}}} \tag{1}$$

Protected subsector of semi-chiral operators, i.e. [Q, O] = 0
In SCFTs: includes those counted by superconformal index

$$\mathcal{I} = \text{Tr}(-1)^F p^{j_1 + j_2 - r/2} q^{j_1 - j_2 - r/2} e^{-\beta \{Q_-, S^-\}}$$
(2)

HOLOMORPHIC TWISTS AND INFINITE SYMMETRIES

- Given any SQFT, we obtain the **Holomorphic Twist** by taking cohomology of any nilpotent supercharge, e.g. *Q* := *Q*_.
 - Anti-holomorphic translations are Q-exact, so twisted theory is (cohomologically) holomorphic [Johansen], [Witten], [Costello]

$$\{Q, \bar{Q}_{\dot{\alpha}}\} = \partial_{\bar{z}^{\dot{\alpha}}} \tag{1}$$

Protected subsector of semi-chiral operators, i.e. [Q, O] = 0
In SCFTs: includes those counted by superconformal index

$$\mathcal{I} = \text{Tr}(-1)^F p^{j_1 + j_2 - r/2} q^{j_1 - j_2 - r/2} e^{-\beta \{Q_-, S^-\}}$$
(2)

Theories are equipped with local product called λ **-Bracket**

$$\{\mathcal{O}_1, \mathcal{O}_2\}_{\lambda} = \oint_{S^3} e^{\lambda \cdot z} d^2 z \ \mathcal{O}_1(z) \ \mathcal{O}_2(0) \tag{3}$$

- Infinite dimensional symmetry enhancements analogous to Virasoro and Kac-Moody, but in 4d [Gwilliam, Williams].
- Higher brackets $\{\mathcal{O}_1, \mathcal{O}_2, \dots, \mathcal{O}_{n+1}\}_{\lambda_1, \lambda_2, \dots, \lambda_n}$

COHOMOLOGIES AND FEYNMAN DIAGRAMS

Feynman diagrams must be Laman graphs.



- Change of variables maps Feynman integrals to a polytope in space of holomorphic loop momenta: the operatope.
- Recursive identities of the operatope used to **bootstrap** Feynman integrals to at least 3-loops (perhaps further!)

COHOMOLOGIES AND FEYNMAN DIAGRAMS

Feynman diagrams must be Laman graphs.



- Change of variables maps Feynman integrals to a polytope in space of holomorphic loop momenta: the operatope.
- Recursive identities of the operatope used to **bootstrap** Feynman integrals to at least 3-loops (perhaps further!)

Polynomials in fields and derivatives ~~ Free Cohomology V

Interacting quantum theory is obtained from underlying free-classical theory V as cohomology of a new operator

$$\mathbf{Q} = Q_0 + Q_1 + Q_2 \dots \tag{4}$$

where Q_n is computed by *n*-loop Feynman diagrams.

All perturbative corrections are contained in the brackets!

$$\mathbf{Q}\mathcal{O} = \{\mathcal{I}, \mathcal{O}\}_0 + \{\mathcal{I}, \mathcal{I}, \mathcal{O}\}_{0,0} + \{\mathcal{I}, \mathcal{I}, \mathcal{I}, \mathcal{O}\}_{0,0,0} + \dots$$
(5)

Holomorphic Confinement of $\mathcal{N}=1$ SYM

• $\mathcal{N} = 1$ SYM is SU(N) gauge theory with an adjoint fermion $\mathcal{L} = \int d^2\theta \frac{-i}{8\pi} \tau \operatorname{tr} W_{\alpha} W^{\alpha} + \text{c.c.}$ (6)

- Twist is identified (in a non-trivial way) with a bc system
- ▶ Free cohomology is gauge invariant polynomials in *b*, *c*, and ∂.

Holomorphic Confinement of $\mathcal{N}=1$ SYM

• $\mathcal{N} = 1$ SYM is SU(N) gauge theory with an adjoint fermion $\mathcal{L} = \int d^2\theta \frac{-i}{8\pi} \tau \operatorname{tr} W_{\alpha} W^{\alpha} + \text{c.c.}$ (6)

- Twist is identified (in a non-trivial way) with a *bc* system
- Free cohomology is gauge invariant polynomials in b, c, and ∂ .
- Adding one-loop corrections changes the cohomology: we find the surviving stress tensor becomes Q-exact.
 - Theory becomes topological at one loop!

Holomorphic Confinement of $\mathcal{N}=1$ SYM

\mathbb{N} $\mathcal{N} = 1$ **SYM** is SU(N) gauge theory with an adjoint fermion

$$\mathcal{L} = \int d^2\theta \frac{-i}{8\pi} \tau \operatorname{tr} W_{\alpha} W^{\alpha} + \text{c.c.}$$
 (6)

- Twist is identified (in a non-trivial way) with a *bc* system
- Free cohomology is gauge invariant polynomials in b, c, and ∂ .
- Adding one-loop corrections changes the cohomology: we find the surviving stress tensor becomes Q-exact.
 - Theory becomes topological at one loop!
- Being topological is compatible with confinement: if topological in the UV, then topological in the IR.
 - Constrains IR physics: the holomorphic twist of the IR must also be topological.



